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COMPENSATION OF SPACE-CHARGE MISMATCH AT TRANSITION OF BOOSTER USING THE TRANSITION-JUMP METHOD

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This method was applied successfully on the CERN PS.

Its application to the NAL accelerators was studied by E. Courant. A theoretical analysis of this method was made by D. Möhl. A more detailed design of such a system for the NAL booster synchrotron is given here based on Möhl's analysis.

Möhl's Analysis

The qualitative and approximate quantitative results of Möhl's analysis are summarized below.

- l. The transition-jump method is effective to partially compensate for the phase-space mismatch caused by space charge at transition which becomes prominent for $\eta_0(0) >> 1$. For best overall matching one needs a delayed jump which starts near transition $(x_1 \sim 0 \text{ in M\"ohl's notation})$. This and $\eta_0(0) >> 1$ are assumed for all subsequent formulas.
- 2. Bunch-length matching for an infinitely fast jump gives a lower limit for the magnitude of the jump

$$\frac{\delta \gamma_{t}}{\delta \gamma} \stackrel{?}{\sim} 1.5 \left(\eta_{o}(0) \right)^{4/3} \tag{1}$$

where $\delta\gamma_t$ is the jump of the transition γ , and $\delta\gamma$ \equiv $\dot{\gamma}$ T (T = time interval about transition in which the adiabatic approximation is



not valid). $\delta \gamma$ is, therefore, the range of γ about transition in which the adiabatic approximation is not valid.

3. For finite jump speed $\left(|\dot{\gamma}_t|\neq\infty\right)$ the increase in phase area A is given by $\left(\frac{\delta A}{A}<<1\right)$

$$\frac{\delta A}{A} \sim \frac{0.60}{(\eta_0(0))^{2/3}} \qquad \frac{\left(\frac{\delta \gamma_t}{\delta \gamma}\right)^2}{\left(1 + \frac{|\dot{\gamma}_t|}{\dot{\gamma}}\right)^{4/3}} \tag{2}$$

which gives the desired jump speed $\dot{\gamma}_t$. If $\frac{\delta\gamma_t}{\delta\gamma}$ is given by Equ. (1) we have

$$\frac{\delta A}{A} \sim 1.4 + \frac{\left(\eta_{o}(0)\right)^{2}}{\left(1 + \frac{\left|\dot{\gamma}_{t}\right|}{\dot{\gamma}}\right)^{4/3}}$$
(3)

4. The growth in $\left(\frac{\Delta p}{p}\right)_{max} \equiv \Delta$ is given by

$$\frac{\Delta}{\Delta_{o}} \sim (\eta_{o}(0))^{-1/3} \left[1 + 0.63 \, \eta_{o}(0) \, \frac{\frac{\delta \gamma_{t}}{\delta \gamma}}{\left(1 + \frac{|\dot{\gamma}_{t}|}{\dot{\gamma}} \right)^{2/3}} \right]$$
 (4)

where Δ_{0} is the $\frac{\Delta p}{p}$ at transition in the absence of space charge.

If $\frac{\delta \gamma_t}{\delta \gamma}$ is given by Equ. (1) we have

$$\frac{\Delta}{\Delta_{o}} \sim (\eta_{o}(0))^{-1/3} + 0.95 \frac{(\eta_{o}(0))^{2}}{\left(1 + \frac{|\dot{\gamma}_{t}|}{\dot{\gamma}}\right)^{2/3}}$$
 (5)

Booster Parameters

The transition parameters $\eta_{_{\scriptsize O}}(0)$ and T can be given in terms of dimensionless parameters by 4

$$\eta_{o}(0) = \frac{\sqrt{3} \pi^{7/2}}{2 (\Gamma(2/3))^{3}} \frac{\mathbf{r}_{p}}{\mathbf{R}} \frac{Ng_{o}}{A^{3/2}} \left(\frac{\mathbf{f}_{rf}}{\dot{\mathbf{r}}}\right)^{1/2}$$
(6)

$$T^{3} = \frac{1}{4\pi} \frac{\gamma^{4}}{f_{rf} \dot{\gamma}^{2} |\cot \phi_{s}|}$$
 (7)

where for the booster at transition the parameters are

 $\Gamma(2/3)$ = factorial function of 2/3 = 1.354

 r_p = classical proton radius = 1.53 x 10^{-18} m

R = machine radius = 75.47 m

N = total number of protons per pulse = 3.5×10^{12}

 g_0 = geometric factor ~ 4.5

 $f_{rf} = rf frequency = 52.20 \times 10^6 sec^{-1}$

 $\dot{\gamma}$ = time rate of increase of γ = 0.407 x 10³ sec⁻¹

 ϕ_s = synchronous phase = 90° ± 20°

A = total phase area occupied by beam in $\frac{p}{mc}$ ϕ

units (A = $\frac{hS}{mcR}$ where S = phase area of each

beam bunch in w ϕ units, and dw = $\frac{R}{h}$ dp) = 0.0069

With these parameters we have for the booster

$$\eta_0(0) = 3.8$$

 $T = 0.28 \times 10^{-3} \text{ sec}$

$$\delta \gamma = \dot{\gamma} T = 0.114$$

and Equ. (1) gives

$$\delta \gamma_{\rm t} \sim 1.0$$
 (8)

For $\frac{|\dot{\gamma}_t|}{\dot{\gamma}}$ = 10 and 20, Equ. (3) gives $\frac{\delta A}{A}$ = 0.83 and 0.35. There-

fore a jump speed of

$$|\dot{\gamma}_{t}| \sim 10^{4} \text{ sec}^{-1}$$

or a jump time for $\delta\gamma_t$ = 1.0 of 0.1 m sec is desirable. This jump speed gives $\frac{|\dot{\gamma}_t|}{\dot{\gamma}} \buildrel 25$ and

$$\frac{\delta A}{A} \sim 0.26, \quad \frac{\Delta}{\Delta_0} \sim 2.2$$
 (10)

γ_t -jump Quadrupoles

Q-jump is a misleading name for the system. What one wants is a γ_t -jump system, and it is quite possible to shift $\gamma_t \text{ by 1 unit without noticeably affecting } \nu_x \text{ and } \nu_y \text{ (or } Q_x \text{ and } Q_y).$ To see this we recall that the equation for the dispersion function x_p after Floquet transformation is

$$\frac{d^2 u}{d\phi^2} + v^2 u = \frac{(v\beta)^{3/2}}{R}$$
 (11)

where

$$u = \frac{x_p}{\sqrt{\nu \beta}}$$

$$d\phi = \frac{dz}{\nu \beta}, \quad z \text{ along orbit}$$

If we add to β a term with zero $\varphi\text{-average}$ (say, $\frac{a}{\nu}$ sin $n\varphi)$, the betatron wave numbers ν_x and ν_y will not be affected. Neglecting other oscillatory terms in β and assuming $\frac{a}{R}$ << 1 we can write approximately

$$\frac{1}{R} (\nu \beta)^{3/2} = \frac{1}{R} (R + a \sin n\phi)^{3/2} = \sqrt{R} + \frac{3}{2} \frac{a}{\sqrt{R}} \sin n\phi$$

The solution of Equ. (11) is then

$$u = \frac{\sqrt{R}}{v^2} + \frac{3}{2} \frac{\frac{a}{\sqrt{R}}}{v^2 - n^2} \sin n\phi$$

The change in orbit length per unit $\frac{\Delta p}{p}$ is

$$\delta L = \oint \frac{x_p}{R} dz = \oint \frac{1}{R} (v\beta)^{3/2} u d\phi$$

$$= \oint \left[\frac{R}{v^2} + \frac{3}{2} a \left(\frac{1}{v^2} + \frac{1}{v^2 - n^2} \right) \sin n\phi + \frac{9}{4} \frac{a^2}{R} \frac{1}{v^2 - n^2} \sin^2 n\phi \right] d\phi$$

$$= \frac{2\pi R}{v^2} \left(1 + \frac{9}{8} \frac{a^2}{R^2} \frac{v^2}{v^2 - n^2} \right)$$

The transition $\boldsymbol{\gamma}_{t}$ is, therefore, changed to

$$\gamma_{t}^{2} = \frac{2\pi R}{\delta L} \stackrel{\circ}{=} \nu^{2} - \frac{9}{8} \frac{a^{2}}{R^{2}} \frac{\nu^{4}}{\nu^{2} - n^{2}}$$
 (12)

If n is chosen to be the integer closest to v, because of the large factor $\frac{v^4}{v^2-n^2}$ a small $\frac{a}{R}$ is adequate to reduce γ_t a great deal.

For the NAL booster the number of cells is N = 24 and ν (namely $\nu_{\rm x}$) is 6.7. We choose n = 6 because it is commensurable with N. Twelve 0.20 m long quadrupoles are placed in the mid-F positions (the cell structure is FOFDOOD) and evenly spaced around the ring. With the 12 quadrupoles excited alternately as F and D, computer runs using SYNCH give for various values of B' of the quadrupoles the following values of $\nu_{\rm x}$, $\nu_{\rm y}$, $\nu_{\rm t}$, and $\nu_{\rm b}$ max at transition.

B' (kG/m)	$\frac{v_x}{x}$	<u>v</u> y	x _{p max} (m)	$\frac{\gamma_t}{}$
0	6.700	6.800	3.189	5.446
±2.825	6.701	6.800	4.157	5.425
±8.476	6.710	6.800	5.966	5.267
±14.127	6.729	6.801	7.512	5.008
±19.777	6.755	6.802	8.728	4.715
±25.428	6.790	6.803	9.600	4.435
±31.079	6.832	6.804	10.156	4.191

We see, therefore, with very modest quadrupole strength one can reduce γ_t by 1 unit. Of course in so doing $x_{p max}$ is increased by a factor of 3. This increase is, of course, unavoidable but it is quite tolerable. The SYNCH output for the cases of B' = 0 and 25.428 kG/m are attached as an appendix.

Acknowledgment

Several interesting discussions with Lloyd Smith were very helpful. The SYNCH runs were made by G. Bellendir.

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